

1 Simulating the Energy Harnessing Capabilities of CdSe 2 Quantum Dots and a Comparison to Traditional Solar Cells

3
4 Rohan Shankar^{1*}, Krish Arora^{1*}, Shrey Birmiwal¹, Shankar Jayaraman²

5
6 ¹ Westwood High School, Austin, Texas

7 ² Marriott International, Baltimore, Maryland

8 **These authors contributed equally to this work*

9

10 Student Authors

11 Rohan Shankar - High School

12 Krish Arora - High School

13 Shrey Birmiwal - High School

14

15

16

17

18

19

20

21

22 SUMMARY

23 Quantum dots are the future of solar energy production, promising conversion efficiencies that
24 might one day rival or even outperform traditional solar panel technology. However,
25 experimental development of quantum dots with the most optimal setups takes months of
26 tedious and costly preparation. To mitigate the cost of this research, simulations of said
27 quantum dots can be used to identify optimal parameters without the need to conduct numerous
28 lengthy experiments. We hypothesize that by taking into consideration the importance of
29 nanoscale quantum phenomena such as confinement and recombination, our developed Python
30 simulation will be able to model the conversion efficiency of CdSe quantum dots in a
31 photovoltaic solar cell with an error of under 5%. Our simulation found CdSe quantum dot
32 conversion efficiency to be 1.66%, which demonstrates comparable rates to the 1.5% efficiency
33 at 50.6% sun found by Lee et. al in 2008 in their experimentation with colloidal CdSe quantum
34 dots— representing a 10.66% realized error and helping prove that simulations can indeed be
35 used as a convincing alternative to painstaking experimentation in the development of optimized
36 quantum dot technologies for energy production. In the future, research should concentrate on
37 reducing this error by factoring in more detailed modeling of quantum phenomenon such as
38 trapping and multiple exciton generation.

39 INTRODUCTION

40 The quest for more efficient energy conversion technologies has spurred considerable interest
41 in the development of advanced photovoltaic materials. Among these, quantum
42 dots—nanoscale semiconductor particles—have emerged as a promising alternative to
43 traditional solar cells due to their potential for high conversion efficiencies and tunable electronic
44 properties. Quantum dots, particularly Cadmium Selenide (CdSe), are noteworthy for their
45 size-dependent electronic and optical characteristics. This phenomenon allows quantum dots to
46 absorb and emit light at specific wavelengths, potentially enhancing their performance in
47 photovoltaic applications. However, the experimental process to optimize these materials is both
48 time-consuming and costly, often involving extensive trial and error.

49

50 Simulations offer a compelling alternative by enabling researchers to model and predict the
51 performance of quantum dots without the need for exhaustive physical experimentation.
52 Leveraging biophysics equations to simulate the electronic properties and energy conversion
53 efficiency of CdSe quantum dots can significantly expedite the development process. The
54 primary challenge is to accurately account for the nanoscale quantum phenomena, such as
55 quantum confinement and recombination dynamics, that influence the efficiency of these
56 materials. Accurate simulation requires incorporating these complex quantum effects into
57 predictive models to achieve reliable and actionable results. Our research aims to address this
58 challenge by developing a Python-based simulation that models the energy conversion
59 efficiency of CdSe quantum dots in photovoltaic cells. We hypothesize that by incorporating
60 detailed quantum mechanical properties and utilizing biophysics equations, our simulation will
61 be capable of predicting the conversion efficiency with an error margin of less than 5%. This
62 hypothesis is grounded in the assumption that accounting for nanoscale quantum phenomena,
63 such as confinement and recombination rates, will yield a simulation result more closely aligned
64 with experimental data.

65

66 The purpose of this study is to validate the accuracy of our simulation model and demonstrate
67 its potential as a tool for optimizing quantum dot-based solar cells. We anticipate that our
68 findings will underscore the viability of simulation as a cost-effective alternative to physical
69 experimentation in the development of advanced photovoltaic technologies. By achieving a
70 simulation accuracy within the proposed error margin, we hope to establish a robust framework
71 for future research and development in quantum dot photovoltaics, ultimately advancing the field
72 towards more efficient and practical solar energy solutions.

73 RESULTS

74 Our simulation showed a predicted solar energy conversion efficiency of 1.66% on average for
75 CdSe quantum dots in our system. When comparing our results to similarly unmodified and
76 unoptimized CdSe quantum dots such as the 50% sun efficiency outlined by
77 <https://pubs.acs.org/doi/full/10.1021/jp802572b> , we get a margin of error of only 10.67%.

78

79 Our first set of results pertains to the simulation of varying solar intensities throughout the year.
80 Solar intensity was adjusted for atmospheric effects and focused by the funnel, recorded hourly.
81 As shown in Figure 1, solar intensity peaked around midday and varied significantly based on
82 the zenith angle. The data illustrates a consistent pattern, with higher intensities observed
83 during the summer months and lower intensities during the winter months. The zenith angle
84 reaches its minimum on day 182, corresponding to the summer solstice. Solar intensity is
85 highest around this day but decreases to even lower than day 91 by the time of day 273,
86 indicating that the period following the summer solstice receives the highest solar intensity,
87 aside from the effect of the clearness index.

88

89 Figure 2 shows the power output, analyzed for checkpoint days 1, 91, 182, and 273. In the line
90 plot shown in Figure 2a, we see that day 182 exhibited the highest overall power generation.
91 However, at hours 10 and 11, day 91 and day 273 showed higher power output, suggesting
92 variations in weather conditions or the probabilistic quantum effects in the quantum dot system.
93 The results indicate that while solar intensity is a significant factor, it is not the sole determinant
94 of power output. The polynomial curve fit in Figure 2b highlights that day 182 consistently had
95 higher power output, but day 91 was comparable at some points, and day 273 showed greater
96 power output than even day 91 despite a lower solar intensity. This is due to the higher
97 clearness index at day 273. This demonstrates the complex interplay of various factors affecting
98 power output, and the importance of considering all of these various factors. Figure 3, showing
99 the cumulative power output throughout the day for different days, showed that the total energy
100 produced increased steadily, with higher slopes around midday corresponding to the more
101 optimal positioning of the sun. This confirms the influence of solar intensity on cumulative
102 energy production.

103

104 Figure 4a provides a visual representation of the temporal and daily variations in power output.
105 It shows that power output peaks between days 169 and 229, corresponding to the summer
106 period with typically sunny days in Austin. Gaps in the heatmap such as those seen on days 94

107 and 253 indicate periods of high cloudiness, which explains their relatively low power outputs
108 despite the days around them having high power outputs. Figure 4b shows the cumulative
109 power output for each day and highlights the influence of the clearness index. Higher cumulative
110 power outputs correspond to higher clearness index values, emphasizing the importance of this
111 metric in power output calculations, but gaps and unpredictability is also seen, which can be
112 attributed to weather conditions.

113

114 Figures 5 and 6 display some of the quantum characteristics of the simulation. Figure 5 shows
115 the size distribution used for simulating the variations in size commonly found in quantum dots.
116 Employing a logarithmic scale was crucial for correctly modeling the quantum dot array, as it
117 reflects the natural distribution of sizes. We used a lognormal distribution with a mean diameter
118 of 2 nanometers and a standard deviation of 0.4 nanometers. Figure 6 shows a sample electron
119 wavefunction. The wavefunction demonstrates the quantum confinement effects in the quantum
120 dot. The oscillatory nature of the wavefunction and the corresponding probability density peaks
121 illustrate the discrete energy levels and the spatial confinement of the electron. The probability
122 density is higher near the center of the quantum dot and decreases towards the edges,
123 indicating that the electron is more likely to be found near the center.

124

125

126 **DISCUSSION**

127

128 The primary results indicate that our simulation effectively models the conversion efficiency of
129 CdSe quantum dots with a realized error of 10.66%. The variations in solar intensity throughout
130 the year, as shown in Figure 1, demonstrate the importance of considering seasonal and daily
131 changes in solar exposure when designing photovoltaic systems. The power output results,
132 especially those depicted in Figures 2 and 3, highlight the complexity of factors influencing
133 energy production, such as weather conditions and quantum effects.

134 The cumulative power output data (Figure 4) emphasize the impact of clearness index and
135 atmospheric conditions on overall power generation. Figures 5 and 6 provide insights into the
136 quantum mechanical behavior of the quantum dots, including the size distribution and
137 wavefunction characteristics, which are crucial for understanding their energy conversion
138 properties.

139 Several factors could have influenced our results. One significant limitation is the accuracy of
140 our atmospheric data, which directly impacts the calculated solar intensity and, consequently,

141 the power output. Additionally, the simplifications made in modeling quantum mechanical
142 processes, such as assuming certain quantum numbers and ignoring some complex
143 interactions, could have introduced errors.

144 The biggest error we have to acknowledge is inaccurate or missing data parameters vital to
145 ensuring the maximum accuracy of our simulation. Since values such as that of the trapping
146 coefficient for CdSe quantum dots are not readily available, we have to manually approximate
147 the value, which introduces significant deviations from the real values.

148 Future research should be aimed at utilizing the most optimal parameters in these simulations,
149 which would ensure that the simulation matches reality much more closely. Moreover,
150 optimization scripts could be incorporated to understand the ideal parameter to use for a
151 quantum dot, which could influence the processes used to create the quantum dot in the
152 experiment.

153 In conclusion, our simulation of CdSe quantum dots demonstrated a conversion efficiency
154 comparable to experimental results, validating the potential of simulation as a tool for optimizing
155 quantum dot technologies. By addressing the limitations and exploring new avenues of
156 research, we can further enhance the efficiency and practicality of quantum dots in solar energy
157 applications.

158

159 MATERIALS AND METHODS

160

161 To accurately simulate the intake of sunlight, we begin by calculating the zenith angle, which is
162 crucial for understanding the interaction of sunlight with our system. The zenith angle, denoted
163 as θ_z , is the angle between the vertical direction and the line to the sun. For our location in
164 Austin, Texas, θ_z is given by:

$$165 \theta_z = \cos^{-1}(\sin(\delta)\sin(\varphi) + \cos(\delta)\cos(\varphi)\cos(H)) \quad (1)$$

166 where δ is the solar declination, φ is the latitude of Austin, and H is the solar hour angle. These
167 parameters vary with time and date and are essential for calculating the precise position of the
168 sun in the sky. δ and H are found from formulae given by NOAA. We then took the cosine of this
169 angle and multiplied it by the solar constant (approximately 1361W/m²) to find the adjusted
170 solar intensity. If $\cos(\theta_z)$ is not positive, this adjusted quantity just becomes equal to 0 as the sun
171 is beneath the horizon at that time.

172 We now need to correct the sunlight intensity value based on the atmosphere. We use the
173 clearness index, for which values will be sourced from NASA's Atmospheric Science Data
174 Center and applied to the adjusted intensity on an hourly basis, ensuring a dynamic adaptation

175 to atmospheric conditions. By multiplying our adjusted intensity with this factor, we get the true
176 value of solar intensity that hits the ground on any given day.

177 The funnel's hyperbolic shape is modeled to focus the incoming solar radiation onto the
178 quantum dot array. The reflective properties of the polished aluminum funnel are crucial here.
179 The concentration factor is calculated based on the funnel's geometry, specifically the ratio of
180 the cross-sectional area of the funnel's wider end to that of the narrower end focused on the
181 quantum dot array. Using this factor and the reflective percentage of polished aluminum, we see
182 that I_{focus} is equal to the product of the reflectivity of polished aluminum (83%), the intensity after
183 atmospheric adjustments, and the concentration factor of 2500 (accounting for the aperture radii
184 of 50 and 1 centimeters, respectively).

185 We calculate the energy per photon using Planck's equation, using the average photon
186 wavelength of 550 nm. This is done by the equation $E_{\text{photon}} = \frac{hc}{\lambda}$. With the total intensity and
187 the energy per photon, we can calculate the number of photons per second striking a area as
188 the product between I_{total} and area divided by E_{photon} . Finally, we obtain the photon generation
189 rate per hour by multiplying the previous value by a factor of 3600. This rate is a crucial input
190 metric for the next part of our simulation, the quantum dot array. To model quantum dots and
191 excitons accurately, we must consider as many major properties and processes as possible,
192 including material composition, geometric structure, quantum mechanical properties, and size
193 dependent properties. This ensures that we calculate the exciton generation rate as accurately
194 as possible. The chosen material for our quantum dots is Cadmium Selenide (CdSe), known for
195 its high absorption efficiency. The distribution of quantum dot sizes is typically characterized by
196 a lognormal distribution, where $D(d)$ is the probability density for diameters, $\mu = \ln(3.9 \times 10^{-9})$
197 represents the mean logarithm of the diameters, and $\sigma = 0.4$ is the standard deviation of the
198 logarithm of the diameters. The distribution will range from 2.3 to 5.5 nanometers. To properly
199 simulate the characteristics of excitons and calculate the rate at which excitons transition
200 between energy states, we numerically solve the Schrödinger equation for a 3D finite potential
201 well. The finite difference method is employed using a grid size of 25 and a step size governed
202 by the radius divided by the grid size. This provides the energy levels of our dot and
203 wavefunctions for electrons and holes. The Laplacian operator ∇^2 of the Schrödinger equation
204 is discretized using the central difference method. General wavefunctions for electrons (ψ_e) and
205 holes (ψ_h) in an ellipsoidal quantum dot with a 3D finite potential well can be approximated as:

$$206 \psi_e(r) \approx \sum_{n,l,m} A_{n,l,m}^e R_{n,l}^e(r) Y_m^l(\theta, \varphi) u_{CB}(r)$$

$$\Psi_h(r) \approx \sum_{n,l,m} A_{n,l,m}^h R_{n,l}^h(r) Y_m^l(\theta, \varphi) u_{VB}(r)$$

208

209 where r denotes the position vector in ellipsoidal coordinates, n, l, m are quantum numbers,
 210 $A_{n,l,m}^{e/h}$ are n, l, m normalization constants, $R_{n,l}^{e/h}(r)$ are the radial functions obtained numerically,
 211 $Y_m^l(\theta, \varphi)$ are spherical harmonics, and $u_{CB/VB}(r)$ are the Bloch functions' periodic parts for the
 212 conduction and valence bands. The normalization constant is calculated by through the
 213 normalization principle, that the total integral of the absolute square of the wavefunction over all
 214 space must always be 1. The radial functions are obtained through the magnitude of the
 215 wavefunction solutions of the Schrödinger equation, calculated numerically by solving the
 216 discretized Schrödinger equation in the finite potential well. The spherical harmonics describe
 217 the angular part of the wavefunction. For simplicity, we assume the quantum numbers l and m
 218 are 0. The Bloch functions describe the periodic part of the wavefunction in the crystal lattice for
 219 the conduction and valence bands, respectively. They are calculated using Bloch's theorem.

220 3.3 Quantum Confinement Effects and Transition Rate

221 Quantum confinement significantly alters the electronic properties of quantum dots. The
 222 effective bandgap energy E_{gap} is modified due to the spatial confinement of charge carriers and
 223 can be expressed as:

$$224 E_{gap} = E_{bulk,gap} + \frac{\hbar^2 \pi^2}{2\mu d^2} - \frac{1.8e^2}{4\pi\epsilon\epsilon_0 d}$$

225 where $E_{bulk,gap}$ is the bulk bandgap energy, d is the quantum dot diameter, and μ is the reduced
 226 mass of the electron-hole pair. The dipole matrix element, essential for calculating transition
 227 rates, is determined by the overlap of the initial and final state wavefunctions. This matrix
 228 element is then used in Fermi's Golden Rule to calculate the transition rate between states. The
 229 rate at which excitons are generated by photon absorption is the product of the photon
 230 generation rate, the quantum efficiency, and the transition rate. These factors collectively
 231 determine the efficiency of exciton creation within the quantum dot system. Exciton
 232 recombination, dissociation, migration, and trapping rates are critical to determining the overall
 233 exciton density. The primary processes considered include recombination rate(The rate at which
 234 excitons recombine typically proportional to the square of the exciton density), dissociation
 235 rate(the rate at which excitons dissociate into free carriers, influenced by the binding energy and
 236 temperature), migration rate(The rate at which excitons move within the quantum dot), trapping
 237 rate(The rate at which excitons are trapped at defects or other localized states), and thermal
 238 dissociation rate(The rate at which excitons dissociate due to thermal energy). The probability of

239 Multiple Exciton Generation (MEG) occurs when the photon energy exceeds a thresh- old,
240 typically 2-3 times the bandgap energy. The rate of Auger recombination, a non-radiative
241 process, is proportional to the cube of the exciton density. Additionally, biexciton formation and
242 recombination are considered, where biexciton formation is influenced by the biexciton binding
243 energy and temperature. The overall exciton density, $n_{exciton}$, considers both generation and loss
244 mechanisms:

$$245 \quad n_{exciton} = \frac{\Gamma_{gen} \cdot (1 + P_{MEG})}{\Gamma_{recomb} + \Gamma_{trap} + \Gamma_{thermal} + \Gamma_{Auger} + \Gamma_{biexciton}}$$

246 where Γ_x represents the rate value of the process x. The net exciton generation rate that
247 reaches the conversion area, accounting for various processes, is calculated as:

$$248 \quad \Gamma_{net} = (\Gamma_{abs} \cdot \eta_{QE} \cdot E_{exciton}) + \Gamma_{dissoc} - \Gamma_{recomb} + \Gamma_{migrate} - \Gamma_{trap} + \Gamma_{thermal}$$

249 This rate, along with the exciton density, directly relates to the overall energy production and
250 efficiency of a quantum dot system. The efficiency and performance of a quantum dot
251 photovoltaic system can be evaluated by calculating the power and energy output. This involves
252 determining the photocurrent generated by excitons, the resultant power output, and the total
253 energy output over a specified time interval. The photocurrent, I_{photo} , is generated by the
254 movement of charge carriers (excitons) in response to light absorption. It is calculated as:

255 $I_{photo} = e \cdot \Gamma_{exciton} \cdot A$, where e is the elementary charge, $\Gamma_{exciton}$ is the exciton generation rate, A is
256 the area of the quantum dot photovoltaic cell. The power output, P_{output} , is the product of the
257 photocurrent and the voltage across the cell. The total energy output, E_{output} , over a given time
258 interval is obtained by multiplying the power output by the duration of the interval.

259

260

261

262

263

264

265

266

267

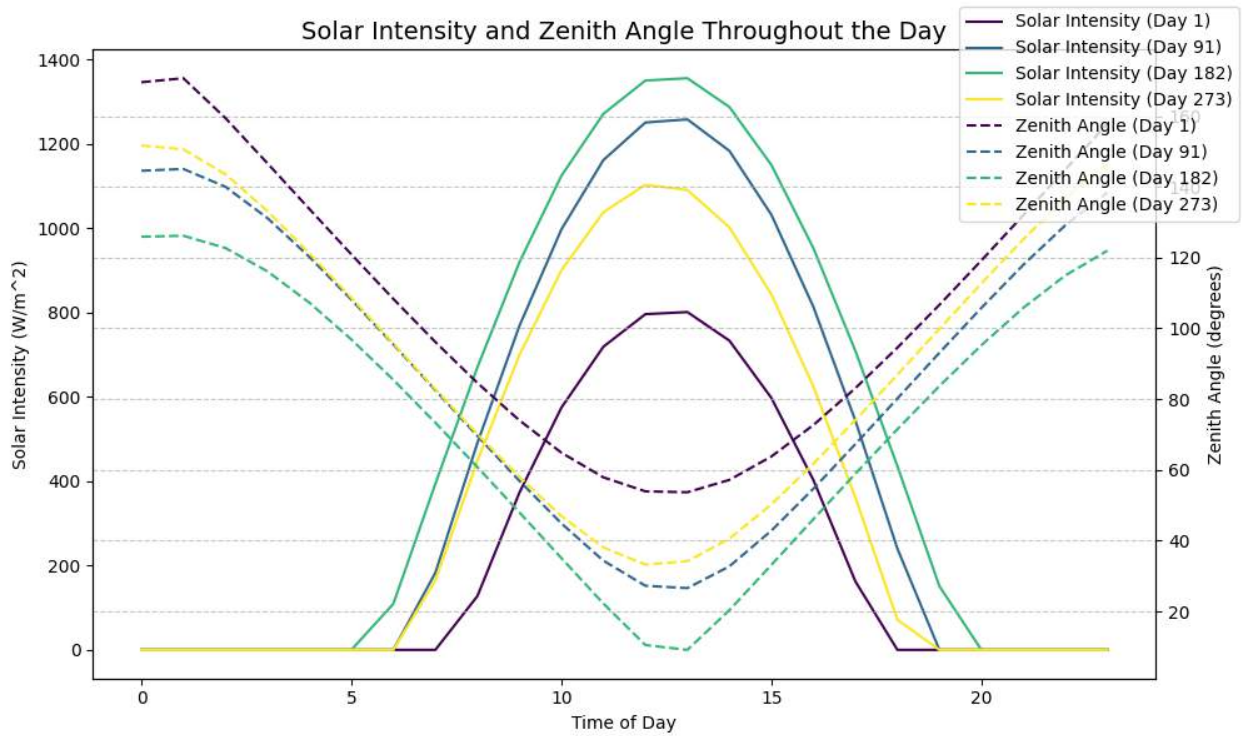
268

269

270

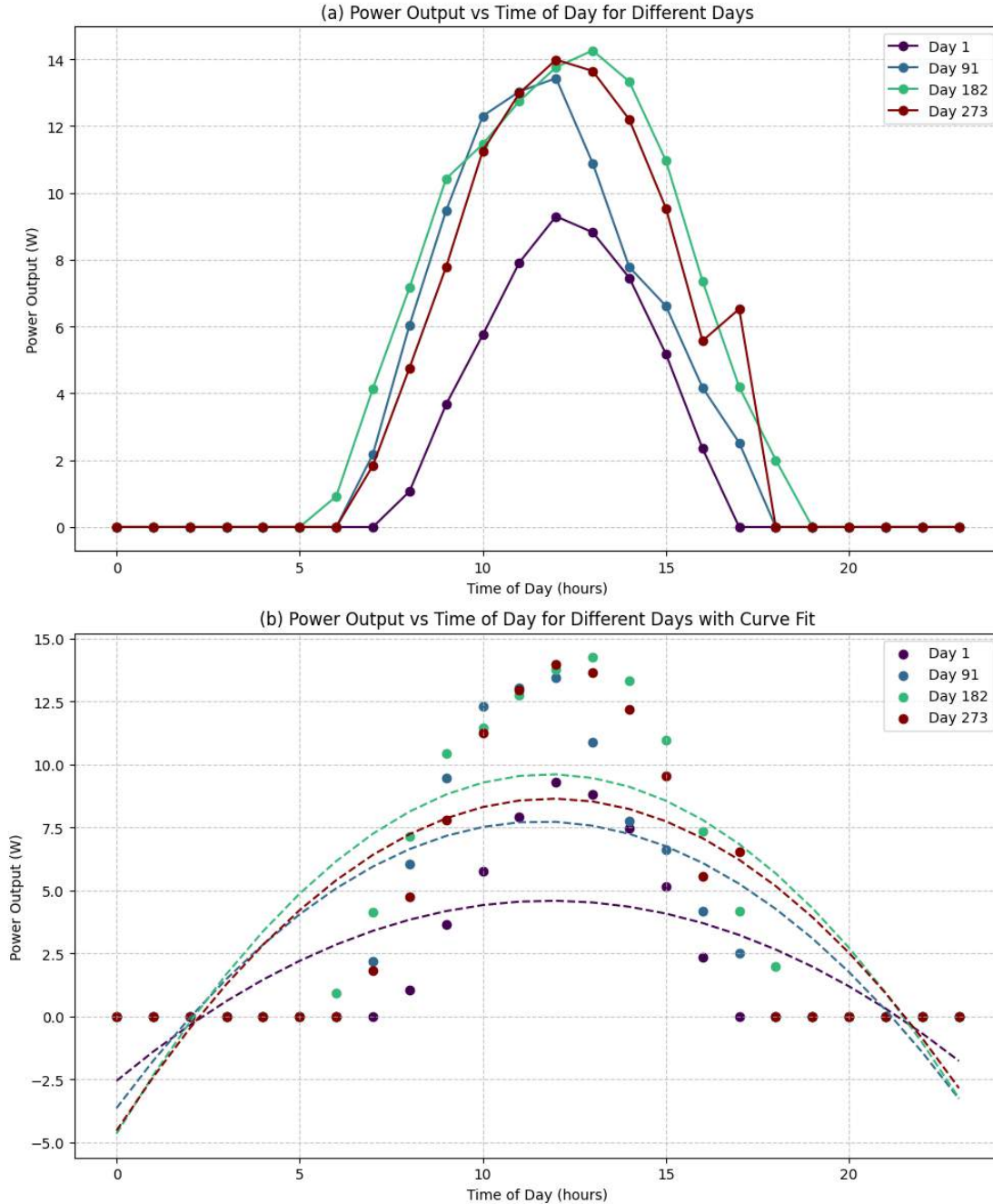
271

272 FIGURES AND FIGURE CAPTIONS



273

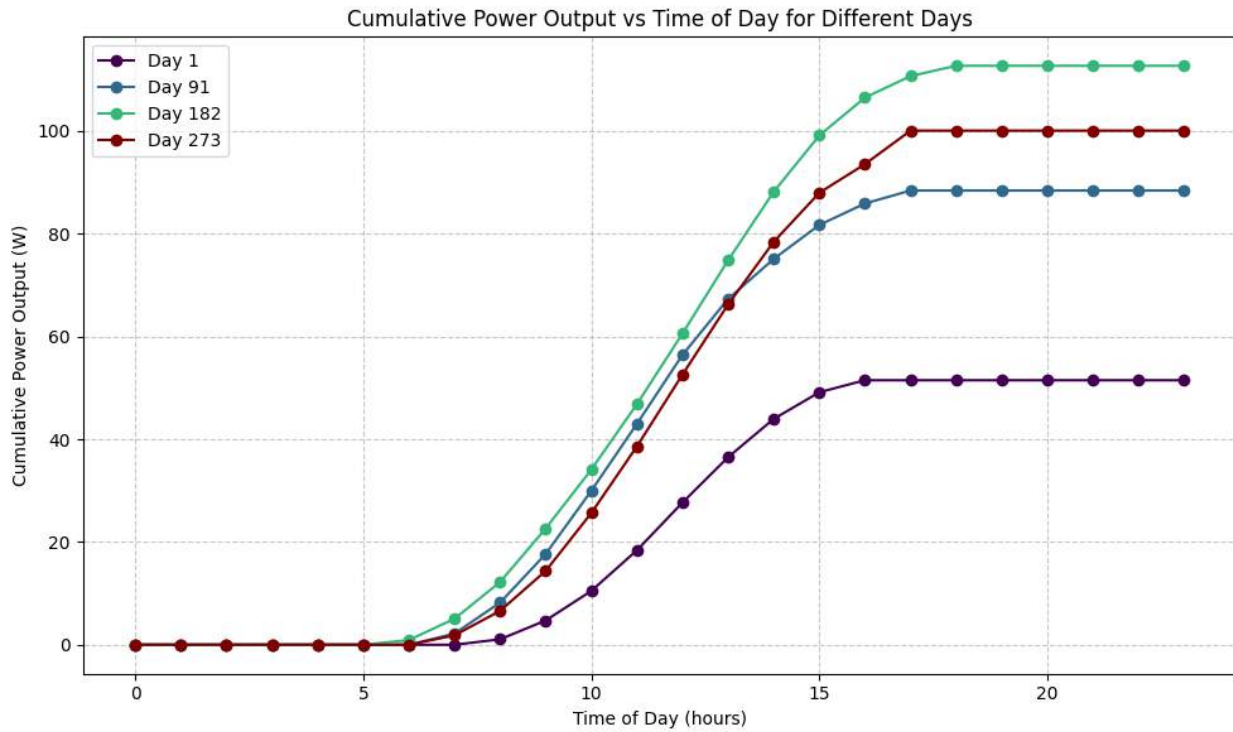
274 Figure 1. Solar Intensity and Zenith Angle Plot Over A Day. Line plot shows the variation in
 275 solar intensity and zenith angle for four days, chosen to represent quarters of the year. The
 276 dotted lines indicate zenith angles and the solid indicate solar intensity.



277

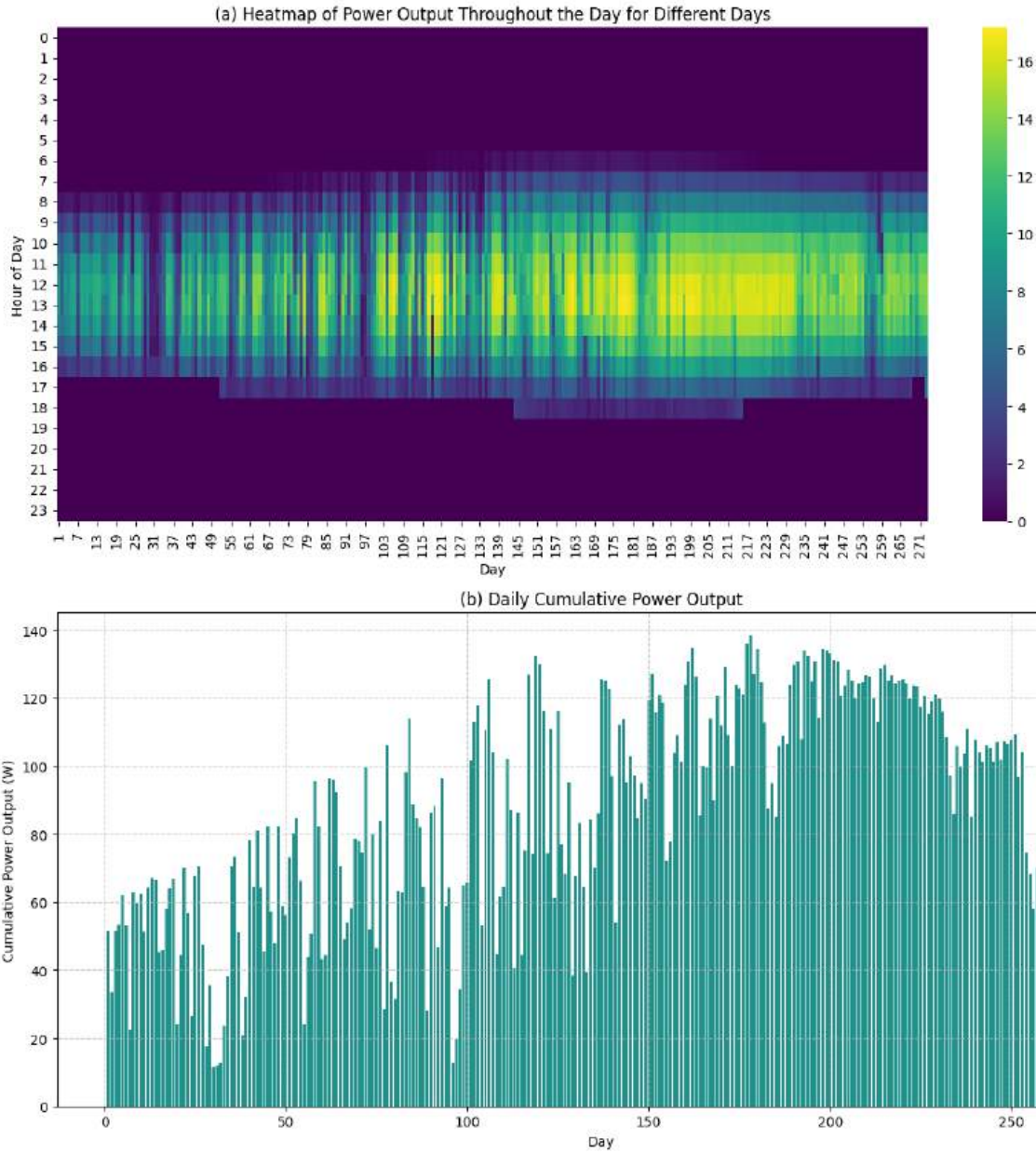
278

279 **Figure 2. Power Output vs Time of Day for Different Days.** (a) This plot shows how the
 280 power output changes over the course of 4 different selected days. It adapts a line fit of the
 281 data, which allows us to see the changes in power of various graphs at different hours. The
 282 individual dots are the true power outputs for a given hour and day. (b) This plot is similar to (a),
 283 but uses a curve fit. This allows us to discern more general trends about how the power output
 284 progresses over a day.



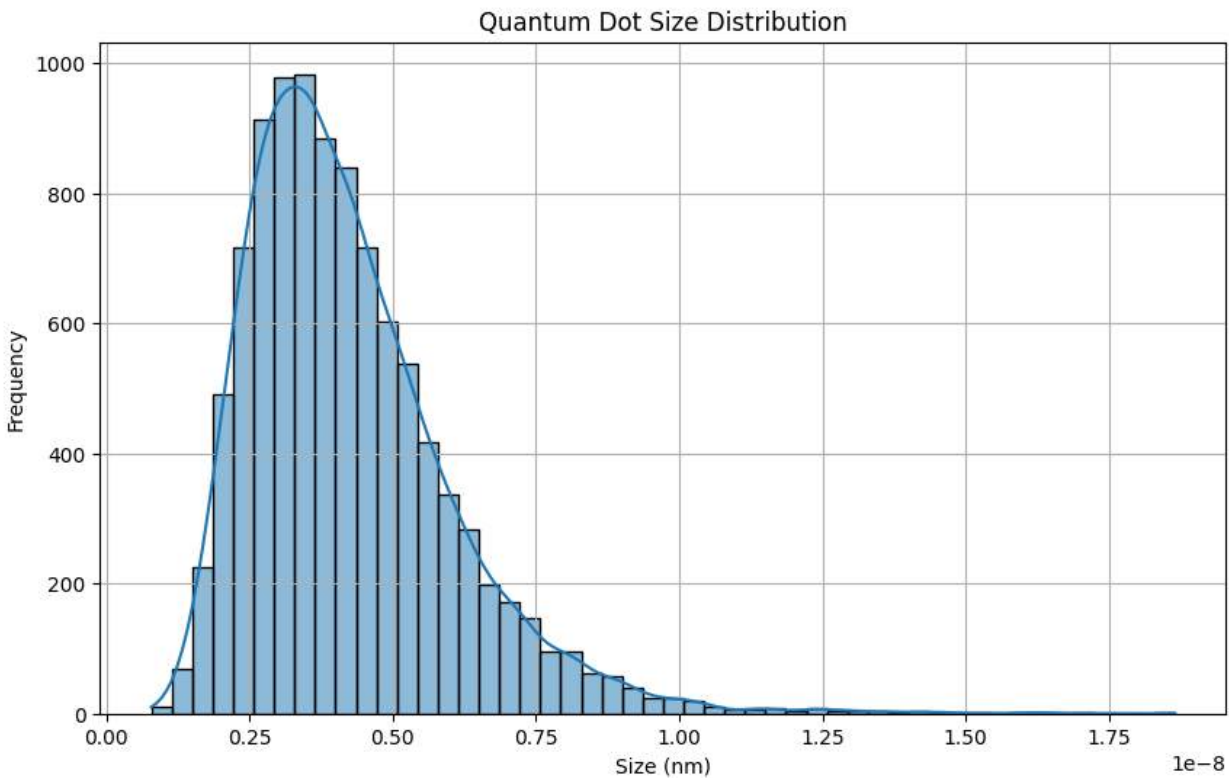
286

287 Figure 3. Cumulative Power Output vs Time for Different Days. Line plot showing the
 288 cumulative power output over the course of four selected days. This plot shows when power
 289 starts being generated, the slopes show when it is at maximum generation, and when power
 290 output stops over the period of a day.



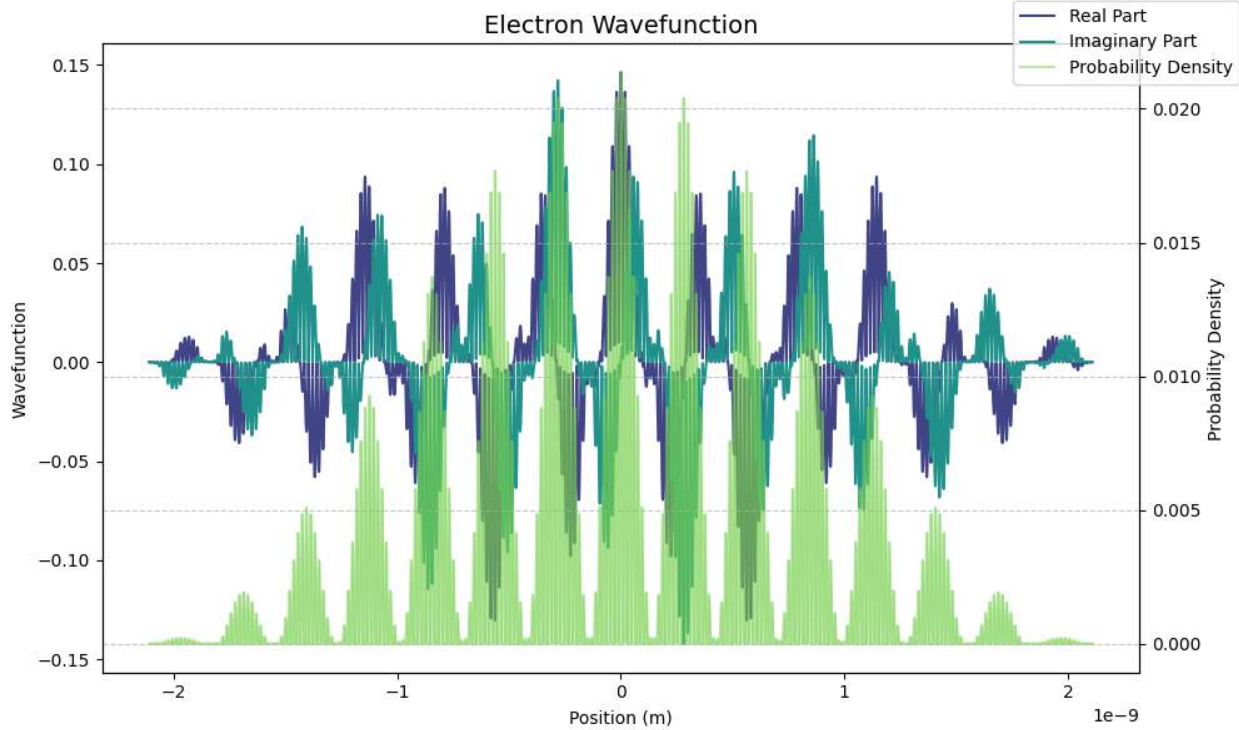
291

292 **Figure 4. Daily Cumulative Power Output.** (a) Heatmap of the predicted power outputs from
 293 January 1st to September 30th. The heat varies over a day, allowing us to see the patterns in
 294 power outputs for a given hour over our date range. (b) This plot shows the cumulative power
 295 output for every day in the year.



296

297 **Figure 5. Quantum Dot Size Distribution.** Shows the distribution of quantum dot sizes in our
 298 simulation. The curve represents the function for CdSe quantum dots, and the frequency
 299 histogram helps identify where the sizes are focused. The distribution is lognormal.



300

301 Figure 6. Example Electron Wavefunction. Shows the real and imaginary parts of the
302 wavefunction, as well as the probability density. The real part of the wavefunction exhibits an
303 oscillatory pattern, characteristic of the particle's quantum state. Nodes, where the wave
304 function crosses zero, indicate points of zero probability density, and the symmetry suggests the
305 potential well's shape and the particle's energy state. The imaginary part complements the real
306 part, providing critical phase information essential for understanding phenomena like
307 interference and tunneling. The probability density, represented by the squared magnitude of the
308 wavefunction, shows the likelihood of finding the particle at different positions. Higher probability
309 density areas indicate where the particle is most likely confined, while the spreading of the
310 wavefunction gives insights into the degree of quantum confinement.

311

312

313

314

315

316

317

318 APPENDIX

319 Code and required dependencies are included in a separate file.